A GMRT Study of Seyfert Galaxies NGC 4235 & NGC 4594: Evidence of Episodic Activity?

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ABSTRACT
Low frequency observations at 325 and 610 MHz have been carried out for two “radio-loud” Seyfert galaxies, NGC 4235 and NGC 4594 (Sombrero galaxy), using the Giant Meterwave Radio Telescope (GMRT). The 610 MHz total intensity and 325−610 MHz spectral index images of NGC 4235 tentatively suggest the presence of a “relic” radio lobe, most likely from a previous episode of AGN activity. This makes NGC 4235 only the second known Seyfert galaxy after Mrk 6 to show signatures of episodic activity. Spitzer and Herschel infrared spectral energy distribution (SED) modelling using the clumpyDREAM code predicts star formation rates (SFR) that are an order of magnitude lower than those required to power the radio lobes in these Seyferts (∼0.13−0.23 M⊙ yr⁻¹ compared to the required SFR of ∼2.0−2.7 M⊙ yr⁻¹ in NGC 4594 and NGC 4235, respectively). This finding along with the detection of parsec and sub-kpc radio jets in both Seyfert galaxies, that are roughly along the same position angles as the radio lobes, strongly support the suggestion that Seyfert lobes are AGN-powered. SED modelling supports the “true” type 2 classification of NGC 4594: this galaxy lacks significant dust obscuration as well as a prominent broad-line region. Between the two Seyfert galaxies, there is an inverse relation between their radio-loudness and Eddington ratio and a direct relation between their Eddington-scaled jet power and bolometric power.

Key words: Seyfert galaxies: general – Seyfert galaxies: individual (NGC 4235, NGC 4594)

1 INTRODUCTION
Seyfert galaxies are one of the several manifestations of the active galactic nuclei (AGN) phenomena, where the primary source of the enormous nuclear energy is mass accretion onto a supermassive black hole. The presence or absence of broad bases to permitted emission lines in the optical spectra has led to the classification of Seyfert type 1s and 2s, respectively (Khachikian & Weedman 1974). Different orientations of the central engine with respect to our line of sight and the resulting obscuration due to a dusty torus, is generally accepted to be the primary explanation for the different Seyfert types, with type 1s being viewed nearly face-on (a.k.a. the Unified Scheme; Antonucci 1993). Seyfert galaxies have traditionally been referred to as “radio-quiet” AGN, i.e., the ratio of their 5 GHz radio flux density to their B-band optical flux density (R), is typically < 10 (Kellermann et al. 1989). However, Ho & Peng (2001) have demonstrated that when the optical nuclear luminosity is extracted through high resolution observations with the Hubble Space Telescope (HST) and the galactic bulge contribution is properly accounted for, then the majority of Seyfert galaxies (belonging to the CfA and Palomar samples in Ho & Peng’s study) transition into the “radio-loud” class. Kharb et al. (2014) confirmed this trend in the Extended 12µm Seyfert sample. Additional criteria (e.g., Panessa et al. 2007) must therefore be adopted to differentiate between the radio emission from Seyferts and radio galaxies that typically exhibit ~100 kpc scale radio jets. Alternatively, more sources that straddle the divide between the “radio-loud” and the “radio-quiet” classes, must be looked at (e.g., Singh et al. 2015b; Mao et al. 2015, and ref-

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Kharb et al. (2006) detected KSRs in (KSRs; Baum et al. 1993; Thean et al. 2000). Gallimore revealed the presence of kiloparsec-scale radio structures continuum in properties between AGN sub-classes. These sources may point towards a genuine AGN: such objects have a weak or missing broad line region (BLR; Panessa & Bassani 2002; Shi et al. 2010). Additionally, Mrk 6 shows signatures of episodic AGN activity.

In order to probe the full extent of the KSRs in two Seyfert galaxies previously known to possess KSRs, viz., NGC 4235 and NGC 4594 (see Gallimore et al. 2006), and to ascertain their origin, we carried out low frequency observations with the Giant Meterwave Radio Telescope (GMRT) at 325 and 610 MHz. The results of this study are presented here. We also present results from an infrared (IR) spectral energy distribution (SED) modelling of NGC 4235 and NGC 4594. Our motivation was to detect any relic emission, if present, at the low radio frequencies of the GMRT, and to examine the relationship between radio lobes and galactic stellar and interstellar medium (ISM) components, which are probed by infrared observations. The ultimate goal is to disentangle the role of the AGN versus the stellar component in producing radio outflows in Seyfert galaxies.

For this paper, we adopt a cosmology where \( H_0 = 73 \text{ km sec}^{-1} \text{ Mpc}^{-1}, \Omega_m = 0.27, \Omega_{\Lambda} = 0.73 \). Spectral index \( \alpha \) is defined such that flux density at a frequency \( \nu \) is \( S_\nu \propto \nu^\alpha \).

## DATA AND ANALYSIS

### 2.1 Sources of Study

NGC 4235 is a nearby (redshift, \( z = 0.008039 \), scale = 0.181 kpc arcsec\(^{-1}\)) Seyfert 1.2 galaxy residing in a nearly edge-on SA(s)a type spiral host with no stellar bar or ring. Seyfert type 1.2 galaxies have ratios of total H\( \alpha \) to [OII] \( \lambda 5007 \) fluxes varying between 2.0 and 5.0 (Winkler 1992; Véron-Cetty & Véron 2006). NGC 4594 is the famous Sombrero galaxy lying at the redshift of \( z = 0.003416 \) (scale =

<table>
<thead>
<tr>
<th>Source</th>
<th>NGC 4594</th>
<th>NGC 4235</th>
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<tbody>
<tr>
<td>Frequency</td>
<td>325 MHz</td>
<td>610 MHz</td>
</tr>
<tr>
<td>Date</td>
<td>06/06/2012</td>
<td>11/06/2012</td>
</tr>
<tr>
<td>Amp Cal</td>
<td>3C 286</td>
<td>3C 147, 3C 286</td>
</tr>
<tr>
<td>( S_\nu ) (Jy)(^{†} )</td>
<td>26.4</td>
<td>36.6, 21.3</td>
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| Phase Cal   | 1311–222, 1311–222 | 1123+055, 1120+143 | 0.090 kpc arcsec\(^{-1}\). It is also an edge-on SA(s)a type spiral galaxy with a LINER\(^{‡} \) nucleus (Ho et al. 1997). It has alternatively been classified as an unobscured “true” type 2 AGN: such objects have a weak or missing broad line region (BLR; Panessa & Bassani 2002; Shi et al. 2010).

### 2.2 Radio Data

Radio observations of NGC 4235 and NGC 4594 at 325 and 610 MHz were carried out in 2012 June (project code 22_066) with the GMRT (Swarup et al. 1991). 30 antennas at bandwidths of 32 MHz at both frequencies were used. A journal of these observations is presented in Table 1. The amplitude calibrators were observed for 20 min at the start and/or end of each run, and the phase calibrators were observed for 5 min every 40 min. NGC 4235 and NGC 4594 were observed for approximately four hours at each frequency.

Data were analysed using the NRAO Astronomical Image Processing System (AIPS). Standard steps in data reduction were followed. The task TVFLG was used to identify and remove the channels affected by radio frequency interference (RFI). Due to lower sensitivity and stability, the channels at the edges of the bands were also removed. Calibration using the primary and secondary calibrators was carried out after editing bad data. The edited and calibrated data on the target source were then averaged in frequency to an extent that kept the bandwidth smearing effect to negligible levels. The data were finally imaged using the wide-field imaging technique, to account for the non-coplanarity of the array. Several iterations of phase-only self-calibration and one iteration of amplitude and phase self-calibration were performed to improve the sensitivity of the images. The final 325 and 610 MHz images presented in Figures 1–3 were created with a ROBUST parameter of 0.25 and 2.0, respectively. The typical rms noise in the 610 MHz images was \( \sim 45 \mu \text{Jy beam}^{-1} \) and in 325 MHz images was \( \sim 150 \mu \text{Jy beam}^{-1} \). We find that the dynamic range of the images at 610 MHz was 1.5 to 3.0 times better than at 325 MHz.

The 325 – 610 MHz spectral index and spectral index noise images were created using the AIPS task COMB, by first obtaining similar resolution images at both frequencies using appropriate weighting with the ROBUST parameter.

\(^{†} \) Flux density of amplitude calibrator set by the SETJY task in AIPS, using the (1999.2) VLA or Reynolds (1934-638) coefficients. \(^{‡} \) Flux density and error for the phase calibrators from AIPS task SETJY.
and subsequently convolving both images with identical circular beams in IMAGR \((8 \times 8 \text{ arcsec}^2)\) for NGC 4235, \(10 \times 10 \text{ arcsec}^2\) for NGC 4594. Images were shifted to be exactly spatially coincident using the task G2EOM. Regions with flux densities below 1.5\(\sigma\) were blanked while using COMB.

### 2.3 Infrared Data

The nuclear structure of galaxies can be probed by infrared data. To this end, we gathered all archival IR data on the two Seyfert galaxies. We then decomposed the infrared spectra and spectral energy distributions using a customised version of the code clumpyDREAM (J. Gallimore, in preparation). clumpyDREAM makes the usual assumption that the infrared SEDs of galaxies can be decomposed into a sum of contributions from stars (photospheric emission), interstellar dust and PAHs (the ISM component), and re-radiation from the dusty torus surrounding an AGN (Genzel & Ceccarelli 2000; Walcher et al. 2011). The nuclear structure of galaxies can be probed by infrared data. To this end, we gathered all archival IR data on the two Seyfert galaxies. We then decomposed the infrared spectra and spectral energy distributions using a customised version of the code clumpyDREAM (J. Gallimore, in preparation). clumpyDREAM makes the usual assumption that the infrared SEDs of galaxies can be decomposed into a sum of contributions from stars (photospheric emission), interstellar dust and PAHs (the ISM component), and re-radiation from the dusty torus surrounding an AGN (Genzel & Ceccarelli 2000; Walcher et al. 2011).

The near-IR data for NGC 4235 come from the 2MASS extended source catalog (Jarrett et al. 2000), while the mid-IR Spitzer IRS\(^4\) data were obtained from Lebouteiller et al. (2011). These data were obtained in staring rather than mapping mode, and so the apertures vary: SL (short-low; \(\sim 5 - 14\) microns) slit = 3.6 arcsec; LL (long-low; \(\sim 14 - 38\) microns) slit = 10.5 arcsec. As a result, any resolved emission will contribute more to the LL module, resulting in a spectral jump at around 14 microns. clumpyDREAM fits for this jump assuming that the AGN components are unresolved but the stars and ISM are resolved. The 60 and 100 micron data were taken from the IRAS\(^5\) point source catalog, and the 250 – 500 micron points were extracted from Herschel SPIRE\(^6\) observations.

The Spitzer IR data for NGC 4594 come from Gallimore et al. (2010). The Spitzer IRS + IRAC\(^7\) spectra were extracted from 20 arcsec synthetic apertures. The near-IR data from 2MASS are aperture-matched to the IRS extraction aperture. The MIPS\(^8\) data were extracted from a 29 arcsec (parallel to slit) \(\times 20\) arcsec (slit width) aperture, assuming a point source geometry. The long-wavelength (250–500 micron) data come from archival Herschel SPIRE observations, extracted as background-subtracted point sources.

The clumpyDREAM code performs a Markov Chain Monte Carlo (MCMC) search over a parameter grid comprising of stellar, ISM and AGN torus components (for more details see appendix in Sales et al. 2015). The code also allows for contributions from very hot (\(\sim 1500\) K), presumably carbonaceous dust grains and AGN-heated dust in the narrow line region. The hot dust component is essentially a part of the dust distribution that the CLUMPY\(^9\) models of Heymann, Nikutta, and Elitzur, do not take into account (Mor et al. 2009). These models assume the same sublimation radius for all grains, when in fact carbonaceous grains have a smaller sublimation radius than silicate grains. Several sources show that the hot dust component is reasonably well modelled by a single blackbody temperature (e.g., Rodríguez-Ardila & Mazzalay 2006).

The torus clouds are distributed in angular altitude \((\theta)\) above the torus equator as \(N(\theta) = N_0 \exp(-\theta^2/\sigma^2)\). From the point-of-view of the AGN or accretion disk, the torus covers a certain fraction of the sky (the "covering fraction" of the torus). Therefore, the integrated luminosity of the torus = \((\text{the AGN luminosity}) \times \text{(the torus covering fraction)}\). The torus luminosity and covering fraction are derived from the model fit, which in turn gives us the AGN bolometric luminosity.

The stellar population is modelled as a single stellar population that undergoes continuous star-formation for 1 Gyr and then passively evolves afterward (Silva et al. 1998). Star formation rate (SFR) is derived using the far-IR relation (equation 4) in Kennicutt (1998). However, since this relation was calibrated for starbursts, there appears to be a \(\approx 30\%\) systematic uncertainty based on comparison to other calibrations. For more quiescent star-forming galaxies, like NGC 4235 and NGC 4594, this calibration may underestimate the SFR by as much as a factor of two. The ISM luminosity is calculated by numerically integrating the IR spectrum of the ISM component from 8 – 1000 microns; it is not bolometric as it does not include UV/visible light that escapes dust extinction.

Upper and lower limits (from larger and smaller apertures) are necessary to provide the best possible constraints on model parameters. Looking at the far-IR region of the SED, fitting the Draine & Li (2007) models strictly to the IRS spectra would result in acceptable models where the predicted far-IR SED would greatly exceed observational limits. There is a similar issue for the contribution of stellar photospheres at the near-IR region of the SED. IRS provides an anchor on starlight only at the shortest wavelengths (~50 microns), but stellar populations present a wide range of near-IR / 5 micron ratios, depending on the mix and age of the underlying populations. Upper limits are necessary to keep that ratio from running away, which then affects how the SED decomposes in the 5-10 micron range, above which dust clearly dominates. Note that we have matched aperture broadband measurements at near-mid IR wavebands (2MASS and IRAC) for NGC 4594. In other words, the fit for NGC 4594 is a fit to 2MASS + IRAC + IRS, but MIPS and Herschel data provide clamps on the far-IR contribution of the Draine & Li models.

### 3 RESULTS

In this Section, we present the results from the radio imaging, followed by those from the IR SED modelling. We first look at some global properties of NGC 4235 and NGC 4594. While Seyfert galaxies have historically been classified as “radio-quiet”, we find that using the total 5 GHz flux density from Gallimore et al. (2006, 11.0 and 5.8 mJy for NGC 4235 and NGC 4594, respectively), and “nuclear” optical B-band (\(4400\) Å) flux density from the HST (0.2 mJy and 72 \(\mu\)Jy for NGC 4235 and NGC 4594, respectively, Papadakis et al. 2008; Maoz et al. 2005), the radio-loudness parameter \((R)\) turns out to be \(\approx 50\) for NGC 4235 and \(\approx 80\)
for NGC 4594, making them both “radio-loud” in the conventional sense. We used the F606W/WFPC2/HST flux density for NGC 4235 and the F330W/ACS/HST flux density for NGC 4594. An optical spectral index of $-0.9$ (e.g., Chibber et al. 1999) was used to convert flux densities between different optical bands.

Dong & De Robertis (2006) have estimated the black hole masses of NGC 4235 and NGC 4594 to be $2.5 \times 10^7 M_\odot$ and $7.9 \times 10^8 M_\odot$, respectively. They used the calibrated black hole - 2MASS K$_s$-band bulge luminosity relation to estimate the same. Using the [O ii] line fluxes listed in Robitaille et al. (2007) and González-Martín et al. (2009) for NGC 4235 and NGC 4594, respectively, and the formulation of Heckman et al. (2004), we have estimated the bolometric luminosities ($L_{bol}$). These are tabulated in Table 3. In Section 4 ahead, we compare these to the more accurate bolometric luminosities obtained from the infrared SED modelling.

### 3.1 Radio Lobes: Morphology, Energetics and Outflow Properties

The double-lobed structure previously observed by Gallimore et al. (2006) is clearly seen in the 325 and 610 MHz images of NGC 4235 (Figure 1). A distinct “hotspot” is seen in the eastern lobe. This was not apparent in the previous VLA observations of Gallimore et al. (2006). The size of individual lobes from the galaxy center is $\approx 12$ kpc; the total radio extent is $\approx 24$ kpc. We also see a hint of diffuse emission at the 3$\sigma$ intensity level, extending beyond the clearly-demarcated western lobe in the 610 MHz image of NGC 4235. However, the higher noise levels in the 325 MHz image prevent us from observing the same (the 3$\sigma$ Jy beam$^{-1}$ level of the feature in our 610 MHz image is $\approx 0.29 \pm 0.19$ in the western lobe and $\approx 0.56 \pm 0.24$ in the eastern lobe.

The double lobes of NGC 4594 are immersed in emission from the galactic disk, which gets more prominent at 325 MHz (Figure 3). The higher frequency (5 GHz) VLA observations of Gallimore et al. (2006) did not detect any galactic emission in this source. However, recent sensitive observations at 6 GHz with the Expanded VLA (EVLA), that achieve image rms noise levels of $\approx 1.5$ mJy beam$^{-1}$, do detect the galactic emission in NGC 4594 (Wiegert et al. 2015). The size of the south-eastern lobe in NGC 4594 from the galaxy center is $\approx 3.5$ kpc, the total radio extent of the lobes is $\approx 5.0$ kpc. These lobes are best seen in Figure 5, where the brightest emission is shown and the disk emission is cut out of the frame. The average 325 – 610 MHz spectral index in the southern lobe is inverted, $\alpha \approx 0.29 \pm 0.16$ (Figure 4). We note that the lobe flux density and spectral index values in NGC 4594 are likely to be contaminated by its galactic emission. The galaxy disk spectral index is on average around $-1.20 \pm 0.25$.

In addition, a “spur” of extent $\approx 3.0$ kpc is observed in the north-eastern direction at both the radio frequencies in NGC 4594 (Figure 5). The spectral index of the “spur” is difficult to extract in the GMRT image due to the presence of the diffuse galactic emission, but appears to be similar to the average spectral index of the lobes. The fact that this feature is not symmetrical about the core and is observed at both the frequencies, gives credence to its reality. The 5 GHz VLA image of Gallimore et al. (2006) did not detect such a feature. Considering that the rms noise achieved at 5 GHz by Gallimore et al. (2006) was $\sim 50$ mJy beam$^{-1}$, and the intensity of the feature in our 610 MHz image is typically around $1$ mJy beam$^{-1}$, a $0.6 - 5.0$ GHz spectral index of $-0.9$ would be implied for this linear feature to fall below 3$\sigma$ at 5 GHz, and therefore not be detected. This spectral index value is not unreasonable for optically thin synchrotron emission from a jet. Interestingly, the 1.5 GHz total and polarised intensity VLA images of NGC 4594 by Bajaja et al. (1988) do show the presence of this feature, although the authors do not comment on it. The “spur” that they refer to is instead the south-eastern lobe/jet that is observed by us and Gallimore et al. (2006). Assuming that the real “spur” intensity is five times the rms noise in figure 2 of Bajaja et al. (1988), we derive a $0.6 - 1.5$ GHz spectral index of $-1.1$, which is again not unreasonable for a jet-like feature. We discuss this feature further in Section 4.

Interestingly, the recent sensitive 1.5 GHz EVLA observations of Wiegert et al. (2015) appear to show a broad northern “plume”, which is much larger in extent than the “spur” and the northern lobe that we observe in the GMRT image, and is at an intermediate position angle (PA) in the sky. The poorer resolution of the EVLA D-array image ($\approx 50$ arcsec) compared to the 610 MHz GMRT image ($\approx 10$ arcsec) does not allow us to make a one-to-one comparison of the various features.

We note that galactic emission was most likely not detected in NGC 4235 at the low radio frequencies of the GMRT, due to its higher redshift compared to NGC 4594. Galactic emission in NGC 4594 that is typically around $350 \mu$Jy beam$^{-1}$ at 610 MHz falls to $6 \mu$Jy beam$^{-1}$ at the redshift of NGC 4235, which is much lower than the noise level at this frequency.

#### 3.1.1 “Relic” Lobe Emission in NGC 4235 (?)

Both the 610 MHz total intensity image and the 325 – 610 MHz spectral index image of NGC 4235 hint at diffuse steep-spectrum radio emission just beyond the well-defined western lobe. This diffuse emission envelopes the western lobe: its shape resembles the shape of the western lobe and it has an extent that is roughly 1.5 times the extent of the western lobe. The average spectral index is $-1.82 \pm 0.17$ in this extended region. This surrounding steep-spectrum emission is clearly reminiscent of “relic” radio emission, as observed for instance, in the lobes of the radio galaxy 3C 388 by Roettiger et al. (1994), where a lobe with an average spectral index of $\sim -0.8$ is surrounded by steep-spectrum emission of spectral index $\sim -1.5$ from a previous AGN activity episode. We note that in principle, a much larger extent of the relic lobe (than 1.5 times the present lobe) cannot be ruled out. However, this extended relic would then have an even steeper spectral index than $-1.8$, making it characteristically different from the relic emission we see, and therefore unlikely.

In order to check if the broken-ring-like region of average 325 – 610 MHz spectral index $-1.8$ could be produced due to the noise peaks in the 325 MHz image, we carried out a simple exercise. We shifted the 325 MHz image suc-
cessively by 25, 35 and 45 pixels in the X-direction, and created 325–610 MHz spectral index images after aligning only pixel (1,1) in both the images, and not the source coordinates. This would superimpose the “relic” emission in the 610 MHz image with noise in the 325 MHz image. We found that where random noise peaks did produce regions of definite spectral index coinciding with the relic emission, the spectral index did vary between $-1.7$ and $-2.0$. However, in no case could we create the broken-ring-like region of average spectral index of $-1.8$ surrounding the western lobe, that is observed in Figure 2. We therefore conclude that while the relic emission is close to the noise level at 325 MHz, its presence is tentatively suggested by its peculiar encompassing morphology. However, by virtue of the 325 MHz relic emission being close to the noise level, the relic spectral index value of $-1.8$, must be regarded as an upper limit.

To the best of our knowledge, the detection of relic lobe emission makes NGC 4235 only the second Seyfert galaxy after Mrk 6 to exhibit multiple episodes of AGN activity. However, relic emission on the hundred-kiloparsec-scale has been observed in the radio galaxy NGC 1167 (Shulevski et al. 2012) and a few other sources at low radio frequencies (e.g., Hurley-Walker et al. 2015; Brienza et al. 2016).

We attempt to obtain the relative ages of the two lobes on the western side of NGC 4235 by using the spectral-ageing formalism of Myers & Spangler (1985). Several assumptions have been made here. We adopt the average spectral index of the eastern lobe ($\alpha \sim -0.6$) for this calculation, and assume that the eastern and western lobes have similar ages. This is a reasonable assumption judging from their similar extents and morphology. The magnetic field is assumed to have the equipartition value of 5 $\mu$G (see Table 2). Furthermore, in order to use figure 2 from Myers & Spangler, it is assumed that the 325–610 MHz and the 1.4–4.9 GHz lobe spectral indices are similar. All these assumptions make the absolute lobe age untrustworthy. We therefore only concern ourselves with the relative ages rather than the absolute ones. Using the electron spectral index $\gamma=2.0$ (or $\alpha=-0.5$) for the Kardashev-Pacholczyk (KP) model, we obtain $X_0=0.22$ and $X_0=0.91$ for $\alpha=-0.6$ for the inner lobe and $\alpha=-1.8$ for the outer “relic” lobe, respectively. This translates to an age of $\sim 40$ Myr for the inner lobe and $\sim 80$ Myr for the outer “relic” lobe. This suggests that the “relic” radio lobe is two times older than the active lobe in NGC 4235.

### 3.1.2 Equipartition Estimates: B-fields and Electron lifetimes

Assuming “equipartition” of energy between relativistic particles and the magnetic field (Burbidge 1959) in the radio lobes, we obtain the minimum pressure, and the particle energy (electrons and protons) at minimum pressure using the relations in O’Dea & Owen (1987). The more sensitive 610 MHz images were used for estimating the equipartition values for individual radio lobes, assuming cylindrical symmetry. We did not make estimations for the short and not-clearly-demarcated north-western lobe of NGC 4594. The total radio luminosity was estimated assuming that the radio spectrum extends from 10 MHz to 100 GHz with a spectral index of $\alpha=-0.6$. Furthermore, it was assumed that the ratio of the relativistic proton to relativistic electron energy was unity. Estimates for two extreme values of the volume filling factor, viz., 1 and $10^{-3}$, were obtained and are presented in Table 2. The extremely small volume filling factor of $10^{-3}$ is justified if we assume that the lobe edges are the projections of precessing jets in Seyfert galaxies (e.g., see Mrk 6 in Kharb et al. 2006). The equipartition magnetic field strength turns out to be $\sim 4 – 30$ $\mu$G in NGC 4235, and $\sim 8 – 60$ $\mu$G in NGC 4594. These parameters are similar to those obtained in other Seyfert galaxies (e.g., Kharb et al. 2006).

The lifetime of electrons in the radio component undergoing both synchrotron radiative and inverse Compton losses on cosmic microwave background (CMB) photons was estimated using the relation by van der Laan & Perola (1969), where the magnetic field was assumed to be the equipartition magnetic field and the break frequency was 1.6 GHz. The magnetic field equivalent to the radiation, which was assumed to be predominantly CMB photons, $B$, was estimated using the relation, $B_r \approx 4 \times 10^{-6} (1+z)^2 \ G$. The electron lifetimes are $\sim 4 – 40$ Myr in the lobes of NGC 4235 and $\sim 2 – 20$ Myr in the southern lobe of NGC 4594. We note again that galactic disk emission contaminates the flux density value for the lobe of NGC 4594.

#### 3.1.3 Can the Radio Lobes be Powered by Stellar Winds?

If we assume that the radio lobe emission has a stellar-wind origin, we can derive the star formation rates required to produce them. We estimated the SFR required to power the lobes using the total lobe luminosity at 610 MHz, and the relations in Condon (1992); Condon et al. (2002) for a Salpeter initial mass function over the mass range 0.1 < $M/M_\odot < 100$. These turn out to be $\sim 2.7 \ M_\odot \ yr^{-1}$ for NGC 4235 and $\sim 2.0 \ M_\odot \ yr^{-1}$ for NGC 4594, for stars more massive than 0.1 $M_\odot$. A radio spectral index of $-0.6$ and +0.3 was assumed for the lobes of NGC 4235 and NGC 4594, respectively, for the SFR estimation.

While these SFR estimates are modest and comparable to those derived in the Milky Way galaxy (e.g., Robitaille & Whitney 2010), they are still an order of magnitude larger than those derived using the Spitzer and Herschel IR data for NGC 4235 and NGC 4594 (Section 3.2). This supports the suggestion that the radio lobe emission in these two Seyfert galaxies is powered by the AGN, rather than by stellar winds. The infrared SFR estimates are from a region only slightly larger than one GMRT beam, and do not include all of the emission from the host galaxy.

Another indirect piece of evidence that favours the suggestion that the radio lobes are AGN-powered, is the “oblique” orientation of the lobes to the host galaxy major axes. Gallimore et al. (2006) have shown that while the lobes can lie at all angles w.r.t. the host galaxy major axes, there appears to be a preference for position angle differences to lie between $40^\circ$ – $60^\circ$. The PA difference is $\sim 40^\circ$ in NGC 4235, and $\sim 50^\circ$ in NGC 4594. Stellar winds, on the other hand, are typically aligned with the galaxy minor axes, where a pressure gradient in the ISM exists (Colbert et al. 1996). This indicates that the radio lobes are “directed” flows, which can more easily be produced by AGN jets rather than nuclear starbursts.
3.1.4 Radio Lobes and Parsec/sub-Kpc Scale Jets

If the Seyfert lobes are AGN-driven, it is essential to look for AGN outflows on parsec-scales in these sources to see if, and how, they align with the kpc-scale radio structures. NGC 4235 has been imaged with VLBI at 8.4 GHz by Anderson et al. (2004). No clear jet-like extension is observed on milliarcsecond-scales. However, the lowest contour levels do indicate a slight extension at a PA of \(\approx 45^\circ\), towards the north-west. The sub-arcsecond-scale VLA A-array 8.4 GHz image of Kukula et al. (1995) clearly reveals a one-sided jet towards the north-west in NGC 4235 at the same PA. The jet in NGC 4235 therefore seems to point towards the northern edge of the north-western lobe. Similar parsec-scale jets and kpc-scale-lobe-edge connections have been observed in other Seyfert galaxies like NGC 6764 and Mrk 6 (Kharb et al. 2006, 2014).

Hada et al. (2013) and Mezcua & Prieto (2014) have observed the parsec-scale jet in NGC 4594 at multiple radio frequencies. A clear S-shaped or wiggly jet with counterjetter is observed in the 5 and 8.4 GHz VLBI images of Hada et al. (2013). The position angle of the parsec-scale jet-counterjet structure is similar to that of the kpc-scale radio lobes in the north-west – south-east direction (Figure 3). Hada et al. (2013) deduce that the northern jet is the approaching one. The parsec-scale jet inclination \((i_j)\) lies in the range \(1^\circ < i_j < 25^\circ\) and the jet speed \((v_j)\) is \(0.02c < v_j < 0.2c\). The parsec-scale jet is therefore sub-relativistic and inclined at a small angle to our line of sight.

The close alignment between the parsec- and sub-kpc-scale radio jets and lobes could in principle suggest that AGN jets are feeding the radio lobes. The non-alignment of the jet and the centres of the lobes (as opposed to the lobe edges) in NGC 4235, could suggest that the jet is curved or S-shaped, as is also observed in Mrk 6 (Kharb et al. 2014). The small angle between the S-shaped jet in NGC 4594 and its torus inclination derived in Section 3.2 (torus inclination \(\sim 17^\circ\)), could further imply that the accretion disk, which is presumably the base of the jet, is warped. Warped accretion disks have been suggested to produce precessing jets (e.g., Pringle 1997).

3.1.5 Lobe Speeds and Kinetic Power

While there is no direct way to measure lobe advance speeds, they can be examined indirectly. Based on the detection statistics of KSRs (\(\geq 44\%\)), which constrain the AGN duty-cycle in Seyferts, Gallimore et al. (2006) have derived relations for the lifetimes and speeds of radio lobes. Using these relations, we derive for \(n = 2\) (two episodes of AGN activity) and a lobe length of 12 kpc in NGC 4235, the duration of a single outflow, \(\tau_{\text{flow}}\), to be \(\sim 20\) Myr, and the lobe speed \(v_{\text{flow}}\), to be \(\sim 600\) km s\(^{-1}\). Interestingly, the lobe advance speed is of the same order as for shock-excited gas inferred from spectroscopic studies of the NLR in Seyfert galaxies (e.g., Bicknell et al. 1998). We note that the outflow duration is about half the age we derived in Section 3.1.1. However, (1) assuming only a slightly different equipartition magnetic field of 7 \(\mu\)G (instead of 5 \(\mu\)G) for the lobe age calculation in Section 3.1.1, gives consistent lobe ages (\(\approx 20\) Myr) from both calculations, (2) the absolute lobe age derived in Section 3.1.1 is unreliable due to the several assumptions made, and (3) the simple relation of Gallimore et al. (2006) assumes that all episodes have identical outflow speeds, which may or may not be true.

Lower limits on outflow speeds can also be derived by dividing the radial extent of the lobes by the electron lifetimes (Baum et al. 1993). The electron lifetimes from synchrotron radiative losses (assuming “equipartition”) and inverse Compton losses due to the CMB background, are presented in Table 2. Therefore, assuming a volume filling factor of unity for the equipartition estimate of the magnetic field, the lower limits on outflow speeds are \(\sim 200 - 250\) km s\(^{-1}\) for the east and west lobes of NGC 4235, and \(\sim 150\) km s\(^{-1}\) for the south lobe of NGC 4594.

We estimate the time-averaged kinetic power \(Q\) of the outflows in the two Seyfert galaxies by using the total 610 MHz flux density of the lobes (without the cores) and the relation from Punsly & Zhang (2011), which has been proposed for powerful radio galaxies. Punsly & Zhang (2011) have slightly modified the relation of Willott et al. (1999) and used an uncertainty parameter \(f\) of 15. The \(f\) parameter includes uncertainty due to minimum energy departures, geometrical effects, filling factors, protonic contribution and the low frequency cutoff. We derive the 151 MHz flux density from the 610 MHz flux density assuming a lobe spectral index of \(\alpha = -0.6\) for NGC 4235 and \(\alpha = +0.3\) for NGC 4594. We obtain a jet kinetic power of \(Q \approx 3.2 \times 10^{40}\) ergs s\(^{-1}\) for NGC 4235 and \(Q \approx 2.1 \times 10^{40}\) ergs s\(^{-1}\) for NGC 4594.

Using the alternate relation \(Q = \frac{E_{\text{min}}}{\eta \times \text{t}_{\text{lobe}} \times \alpha}\) (Rawlings & Saunders 1991), we obtain the same \(Q\) estimate (\(= 3.2 \times 10^{40}\) ergs s\(^{-1}\)) for NGC 4235. For this we have used the minimum energy, \(E_{\text{min}} = 2 \times 10^{52}\) ergs (Table 2; a representative value was chosen for an intermediate filling factor between 1 and \(10^{-3}\)), lobe age = 40 Myr (see Section 3.1.1) and \(\eta = 0.5\) (parameter that allows for work done on external medium; Rawlings & Saunders 1991). Kinetic powers of the order of \(\sim 10^{40} - 10^{41}\) ergs s\(^{-1}\) have been found in other low-luminosity AGN with direct estimates of mass outflow rates (e.g., NGC 1433, Combes et al. 2013). This gives us confidence that the relations for time-averaged kinetic power derived for powerful radio galaxies, also seem to work in Seyfert galaxies with radio lobes.

3.2 Nuclear Properties from the IR SED

Modelling

The infrared SED modelling with the customised clumpy-DREAM code decomposes the stellar and AGN contribution to the nuclear emission in NGC 4235 and NGC 4594. Modelling reveals that the best-fit model for NGC 4594 includes (i) a clumpy dusty torus, (ii) diffuse ISM, and (iii) a stellar population. The best-fit model for NGC 4235, on the other hand, includes these three components plus an additional, presumably carbonaceous, hot dust component with temperature \(\sim 1500\) K, that is described in Section 2.3. The hot dust component is found only in quasi-stellar objects (QSOs) and type 1 AGN. The SED best-fit models are presented in Figures 5, 6 and Tables 4, 5. Estimators of the probability densities of the fitted parameters for both Seyfert galaxies are shown in Figures A1-A4. The model parameters are described in the caption to Table 4. The jump in slit width between the short-low SL and the long-low LL modules mentioned in Section 2.3, can be seen in the model.
spectral index of $\lambda$IR = 42.30 to 42.99 for NGC 4235 and 43.385 ± 0.012 for NGC 4594. Since NGC 4235 does not have 12 and 25 micron detections but only upper limits, the values for luminosity and star formation rates are ranges rather than measured values with uncertainty. Using the relations in Kennicutt (1998), we obtain global SFR values of 0.089 − 0.438 $M_\odot$ yr$^{-1}$ for NGC 4235 and 1.09 ± 0.30 $M_\odot$ yr$^{-1}$ for NGC 4594. Keeping in mind the factor-of-two systematic uncertainty mentioned earlier in Section 2.3, the upper SFR limit for NGC 4235 is 0.88 $M_\odot$ yr$^{-1}$. Overall, these global values are greater than the nuclear values estimated for NGC 4594. As discussed ahead in Section 4, these estimates are more accurate than the ones for NGC 4235 and NGC 4594, making them both type 1 (in the sense of being unobscured) in the Unified scheme.

The bolometric AGN luminosity derived from the IR SED modelling, which relies on the torus luminosity and its covering fraction, is $3.5 \times 10^{42}$ ergs s$^{-1}$ for NGC 4235 and $1.4 \times 10^{42} \pm 0.012$ ergs s$^{-1}$ for NGC 4594. As discussed ahead in Section 4, these estimates are more accurate than the ones derived from the [O III] line fluxes in Section 3, and are used for estimating relevant parameters like the Eddington ratio ($\lambda_{\text{Edd}} \equiv L_\text{bol}/L_\text{Edd}$), where the Eddington luminosity is $L_\text{Edd} = 1.25 \times 10^{38} M_\odot$ yr$^{-1}$, and kinetic to bolometric power (or jet to accretion efficiency); see Table 3).

3.2.1 Global Star Formation Rates from IRAS and Radio Data

We can derive “global” star formation rates using infrared fluxes from IRAS. We estimated the global IR luminosity, $L(\lambda = 8$ − 1000 microns), using Sanders & Mirabel (1996), and subtracted off the AGN $L(\lambda = 8$ − 1000 microns) component based on the SED model fit. The AGN-subtracted, IRAS estimates of the global IR luminosity in ergs s$^{-1}$ are: log $L_{\text{IR}} = 42.30$ to 42.99 for NGC 4235 and 43.385 ± 0.012 for NGC 4594. Since NGC 4235 does not have 12 and 25 micron detections but only upper limits, the values for luminosity and star formation rates are ranges rather than measured values with uncertainty. Using the relations in Kennicutt (1998), we obtain global SFR values of 0.089 − 0.438 $M_\odot$ yr$^{-1}$ for NGC 4235 and 1.09 ± 0.30 $M_\odot$ yr$^{-1}$ for NGC 4594. Keeping in mind the factor-of-two systematic uncertainty mentioned earlier in Section 2.3, the upper SFR limit for NGC 4235 is 0.88 $M_\odot$ yr$^{-1}$. Overall, these global values are greater than the nuclear values estimated in Section 3.2. This is not surprising since these galaxies are disky and the global values include all of the disk.

Using the 610 MHz flux density from the galactic disk in NGC 4594 (derived by subtracting the core-lobe flux density from the total observed), we derive an SFR of $\sim 0.82 M_\odot$ yr$^{-1}$ for stars more massive than 0.1 $M_\odot$. (A spectral index of −1.2 was assumed for the disk emission; see Section 3.1.2.) This radio estimate is consistent with the global SFR from the IRAS data, given the uncertainties. A similar analysis cannot be carried out for NGC 4235 because galactic disk emission is not detected at either of the GMRT frequencies. We conclude that the global SFR estimates are consistent, whether using radio or infrared fluxes. This strongly suggests that the difference in the “nuclear” SFR estimates between the infrared and the radio, is significant.

4 DISCUSSION

Observations of the Seyfert galaxies NGC 4235 and NGC 4594, at the low radio frequencies of 325 and 610 MHz, tentatively reveal signatures of episodic AGN activity in them. Both the 610 MHz total intensity and the 325−610 MHz spectral index images suggest the presence of a “relic” radio lobe in the Seyfert galaxy NGC 4235. Based on a simple spectral ageing analysis, the relic outer lobe appears to be at least two times older than the present lobe. This implies that the AGN in NGC 4235 was switched “off” for the same time that it has been “on” for the current episode. If we assume that the lobe expansion speed was constant for the two activity episodes, the outer lobe should have been at least two times in overall extent than the inner lobe. However, this is not the case (extent is about 1.5, see Section 3.1 and Figure 1), we can conclude that the lobe expansion was faster in the inner lobe compared to the outer one. This is plausible if the previous AGN activity episode cleared out the surrounding medium, so that the new lobe propagates in a rarer medium. The detection of the “relic” lobe makes NGC 4235 only the second known Seyfert galaxy after Mrk 6 to exhibit lobe emission from multiple AGN activity episodes.

A $\sim 3$ kpc linear “spur-like” feature is observed nearly perpendicular to the double-lobed structure in NGC 4594, at 325 and 610 MHz. This feature was not detected in the 5 GHz VLA image of Gallimore et al. (2006), but appears to have been detected in the 1.5 GHz VLA image of Bajaja et al. (1988), although the authors do not comment on it. Both these observations indicate a spectral index of around $−1.0$ between 610 MHz and the GHz frequencies. It is interesting to note that the Milky Way Spur, which is $\approx 4.0$ kpc long and extends perpendicular to the Galactic plane, has been suggested to be a one-sided jet or a magnetic tornado produced by a twisted poloidal magnetic field between the disk and halo (Sofue et al. 1989). Could the “spur” in NGC 4594 have a similar origin ? Since the VLBI jet in NGC 4594 is perpendicular to this linear feature, the AGN does not seem to be currently fuelling it. Since the new data of Wiegert et al. (2015) show a much larger northern “plume” along the galaxy minor axis at a PA that is intermediate between the “spur” PA and the lobe PA, we wonder if the “spur” is part of a larger relic lobe emission in NGC 4594. However, as of now, the origin of this peculiar feature remains unclear.

Constraints on the star formation rates derived through infrared observations, support the suggestion that the radio lobes are AGN- rather than stellar-wind-powered, in both the Seyfert galaxies. The best supporting evidence for the AGN-connection comes from the detection of parsec-scale radio jets in both Seyfert galaxies, with position angles sim-
ilar to the kpc-scale lobe emission, indicative of a close connection. However, jet curvature, going from parsec to kiloparsec scales, is also implied. Curved jets could be indicative of precessing accretion disks, or instabilities in the radio outflows (e.g., Pringle 1997; Hardee 1982). Kharb et al. (2006) had modelled the entire radio lobe in Mrk 6 to be the projection of a precessing jet. Edge-brightened lobe-like structures (or figure-of-eight-shaped structures) can indeed be produced in hydrodynamical jet simulations of slowly precessing jets (M. Nawaz and G. Bicknell, private communication). Based on the torus inclination derived from the SED modelling of NGC 4594, the radio jet is not perpendicular to the torus; rather it lies between \( \approx 8^\circ \) and \( 16^\circ \) to the torus axis. This could suggest a warp in the accretion disk, which could in turn explain the S-shaped parsec-scale jet as a precessing jet emerging from the accretion disk in NGC 4594.

AGN bolometric luminosities have been estimated using the [O\textsc{iii}] line fluxes, as well as the IR SED modelling. These are tabulated in Table 3. The \( L_{\text{bol}} \) values derived from the [O\textsc{iii}] line fluxes are eight (for NGC 4235) and four (for NGC 4594) times larger than the \( L_{\text{bol}} \) values derived from the infrared SED. This is likely to be due to the large uncertainties associated with the bolometric calibration from [O\textsc{iii}] as mentioned in Heckman et al. (2004): the derived bolometric luminosities may be off by as large a factor as 14 (using 3\( \sigma \) quoted in Heckman et al.). We have therefore used the \( L_{\text{bol}} \) values from SED modelling for all calculations. We obtain Eddingtion ratios of 0.001 for NGC 4235 and \( \sim 1 \times 10^{-5} \) for NGC 4594. While the Eddingtion ratios of NGC 4235 and NGC 4594 are typical of those derived in Seyfert galaxies (e.g., Ho 2008), it is interesting to note that between the two sources, the source with the higher radio-loudness parameter, \( \nu R \), NGC 4594, has the smaller Eddingtion ratio. Such a trend has indeed been observed for low luminosity AGN (e.g., Panessa et al. 2007; Merloni & Heinz 2007). Rafter et al. (2009) have identified a large number of broad-line AGN having radio-loudness parameters between 10 and 100, and referred to them as “radio-intermediate”. NGC 4235 and NGC 4594 fall at the extreme left edge (or beyond) of their log \( R \) versus log \( (\lambda_{\text{Edd}}) \) plot (their figure 7), just above the log \( R = 1 \) “radio-loudness” dividing line. They are both clearly “radio-intermediate” by the definition of Rafter et al. (2009) and extend the log \( (R) - \) log \( (\lambda_{\text{Edd}}) \) analysis to much smaller Eddington ratios.

The kinetic to bolometric power \( (Q/L_{\text{bol}}) \) is representative of the jet to accretion efficiency: it is \( \sim 0.009 \) for NGC 4235 and \( \sim 0.002 \) for NGC 4594. We list the X-ray bolometric correction factors \( (L_{\text{x}}/L(2-10 \text{ keV})) \), the Eddington-scaled bolometric powers \( (L_{2-10 \text{ keV}}/L_{\text{Edd}}) \), and jet powers \( (Q/L_{\text{Edd}}) \) for the two sources in Table 3. The absorption-corrected 2 – 10 keV X-ray luminosity \( L(2-10 \text{ keV}) \) was obtained from Liu et al. (2014): it is \( 4.0 \times 10^{44} \) ergs s\(^{-1}\) in NGC 4235 and \( 3.8 \times 10^{40} \) ergs s\(^{-1}\) in NGC 4594. The X-ray bolometric correction factors tell us how much the X-ray intensity has to be scaled up to obtain the total accretion energy density (Vasudevan & Fabian 2007). These range from \( \sim 10 \) to \( \sim 40 \) for the two galaxies.

We find that the two Seyfert galaxies follow the same upward linear trend that is observed in the \( (L_{2-10 \text{ keV}}/L_{\text{Edd}}) \) plot of Merloni & Heinz (2007, figure 2), but are not coincident with their best-fit line, even after the X-ray luminosity is multiplied by the bolometric correction of 5 chosen by them. Since all the relevant estimates in this analysis are obtained from the literature, it is difficult for us to identify the reason for, or the significance if any, of this offset. Overall, for NGC 4235 and NGC 4594, the bolometric power exceeds the jet power, and is linearly related to it. This implies that the optical AGN is capable of driving the radio outflow.

Ho (2008) found out that the BLR appears to vanish at the lowest Eddington ratios. This seems to be true in the case of NGC 4594, which has a very small Eddington ratio of \( 5 \times 10^{-5} \) for a black hole mass of \( 7.9 \times 10^{8} \) M\(_{\odot}\) (Dong & De Robertis 2006). The IR SED modelling indicates that the extinction due to a dusty torus is not significant in NGC 4594, as its inclination is nearly face-on. Additionally, the escape probability of photons from the accretion disk/BLR is essentially 100%. Our data are therefore consistent with NGC 4594 being a “true” type 2 AGN with no torus extinction and no BLR.

5 SUMMARY AND CONCLUSIONS

Low radio frequency observations with the GMRT reveal unprecedented details in the lobes and galactic disk of two nearby “radio-loud” Seyfert galaxies, NGC 4235 and NGC 4594.

(i) The 610 MHz total intensity image and the 325 – 610 MHz spectral index image reveal tentative evidence of a “relic” radio lobe in NGC 4235, pointing to episodic AGN activity in this galaxy. This relic outer lobe is at least two times older than the present lobe. A \( \sim 3 \) kpc linear, steep-spectrum “spur-like” feature of unknown origin, is observed nearly perpendicular to the double-lobed structure in NGC 4594. Galactic disk emission is clearly detected at both radio frequencies in NGC 4594, but not in NGC 4235. Galactic contamination makes the flux density and spectral index values for the lobe in NGC 4594 uncertain.

(ii) The infrared SED modelling suggests that the nuclear star formation rate is 0.23 M\(_{\odot}\) yr\(^{-1}\) and 13 M\(_{\odot}\) yr\(^{-1}\) in NGC 4235 and NGC 4594, while the SFR required to power the radio lobes entirely is 2.7 M\(_{\odot}\) yr\(^{-1}\) and 2.0 M\(_{\odot}\) yr\(^{-1}\), respectively. This supports the idea that the lobes are AGN, rather than stellar-wind powered. The global SFR estimates from both the infrared and radio data are consistent within the errors for the galactic disk emission in NGC 4594.

(iii) The lobe-AGN connection is further strengthened by the presence of parsec-scale jets in NGC 4594 and a sub-kiloparsec-scale radio jet in NGC 4235, that are roughly at the same position angle as the radio lobes, and are therefore likely to be powering them. The SED model-fitting and VLBI imaging suggests that the parsec-scale jet in NGC 4594 is inclined at \( \approx 8^\circ \) to 16\(^\circ\) to the torus axis: its S-shape could therefore indicate a precessing jet emerging from a warped accretion disk. A curved jet is also implied in NGC 4235.

(iv) The SED modelling strongly supports the suggestion that NGC 4594 is a “true” type 2 AGN with insignificant dust obscuration and a weak or absent BLR. The dusty torus inclination is 4\(^\circ\) in NGC 4235 and 17\(^\circ\) in NGC 4594, while the “escape probability” for the accretion disk and BLR photons to escape the torus, is essentially 100% in both NGC 4235 and NGC 4594, making them both type 1 (in the sense of...
being unobscured) in the Unified scheme. The SED fitting reveals that NGC 4235 has an additional hot dust component (temperature $\sim 1500$ K), apart from the torus, diffuse ISM and stellar components.

(v) AGN bolometric luminosities derived from the [O III] line fluxes are four to eight times larger than the more accurate luminosities derived from the IR SED modelling. The Eddington ratios obtained from the latter are $\sim0.001$ for NGC 4235 and $\sim 1 \times 10^{-5}$ for NGC 4594.

(vi) The time-averaged kinetic power of the outflow is $3.2 \times 10^{40}$ ergs s$^{-1}$ in NGC 4235. Since galactic emission contaminates the total lobe emission in NGC 4594, an upper limit of the kinetic power is $2.1 \times 10^{40}$ ergs s$^{-1}$. Such kinetic powers have been found in other low-luminosity AGN with direct estimates of mass outflow rates. The relations for kinetic power that are derived for powerful radio galaxies are also valid for Seyfert galaxies with radio lobes. The jet to accretion efficiency factor, represented by $Q/\dot{L}_{\text{bol}}$, is $\sim0.009$ for NGC 4235 and $\sim0.002$ for NGC 4594. The X-ray bolometric correction factors range from $\sim10$ to $\sim40$ for the two galaxies.

(vii) Between the two Seyfert galaxies, there is an inverse relation between the radio-loudness and Eddington ratio and a direct relation between the Eddington-scaled jet power and bolometric power. Such trends have indeed been found for low luminosity AGN in the literature.

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NGC4235 and NGC4594 with the GMRT

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Kharb et al.

Wiegert T., et al., 2015, AJ, 150, 81
Table 2. Equipartition Estimates for the Lobes

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<th>Comp</th>
<th>$S_{tot}$</th>
<th>Major</th>
<th>Minor</th>
<th>$\phi$</th>
<th>$L_{rad}$</th>
<th>$B_{min}$</th>
<th>$E_{min}$</th>
<th>$P_{min}$</th>
<th>$E_{tot}$</th>
<th>$U_{tot}$</th>
<th>$t_e$</th>
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<tr>
<td></td>
<td>mJy</td>
<td>arcsec</td>
<td>arcsec</td>
<td>rad</td>
<td>$\times 10^{39}$</td>
<td>$\times 10^{-6}$</td>
<td>$\times 10^{54}$</td>
<td>$\times 10^{-12}$</td>
<td>$\times 10^{54}$</td>
<td>$\times 10^{-11}$</td>
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</tr>
<tr>
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<td>8.6</td>
<td>52.1</td>
<td>47.2</td>
<td>1</td>
<td>15.8</td>
<td>4.5</td>
<td>65.8</td>
<td>1.9</td>
<td>82.3</td>
<td>0.4</td>
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<td>40.1</td>
<td>1</td>
<td>7.8</td>
<td>4.3</td>
<td>35.2</td>
<td>1.7</td>
<td>44.0</td>
<td>0.4</td>
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<td>7.9</td>
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<td>5.8</td>
<td>4.2</td>
<td>1.2</td>
<td>23.1</td>
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</table>

Images at 610 MHz were used for all estimates. A lobe spectral index of $-0.6$ was adopted throughout. Column 1: Source name followed by lobe (w=west, e=east, s=south). Column 2: Integrated flux density. Columns 3 & 4: Major and Minor axes in arcseconds, respectively. Column 5: Volume filling factor. Column 6: Total radio luminosity. Column 7: Minimum magnetic field strength. Column 8: Minimum energy. Column 9: Minimum pressure. Column 10: Total energy in particles and fields $E_{tot} = 1.25 \times E_{min}$. Column 11: Total energy density $U_{tot} = E_{tot} (\phi V)^{-1}$. Column 12: Electron lifetimes in years.

Table 3. Eddington-scaled Quantities and Star-formation rates

<table>
<thead>
<tr>
<th>Source</th>
<th>$M_{bh}$/$M_\odot$</th>
<th>$L_{bol}$/$L_{Edd}$</th>
<th>$L_{Edd}$/$L_{bol}$</th>
<th>$\lambda_{Edd}$</th>
<th>$L_{bol}/L_{2-10~\text{keV}}$</th>
<th>$L_{2-10~\text{keV}}/L_{Edd}$</th>
<th>$Q/L_{Edd}$</th>
<th>SFR-1/$M_\odot$ yr$^{-1}$</th>
<th>SFR-2/$M_\odot$ yr$^{-1}$</th>
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<tr>
<td>NGC 4235</td>
<td>$2.5 \times 10^7$</td>
<td>$2.9 \times 10^{43}$</td>
<td>$3.5 \times 10^{42}$</td>
<td>$3.1 \times 10^{45}$</td>
<td>$0.001$</td>
<td>$9$</td>
<td>$0.0001$</td>
<td>$1.0 \times 10^{-5}$</td>
<td>$2.7$</td>
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<tr>
<td>NGC 4594</td>
<td>$7.9 \times 10^8$</td>
<td>$5.2 \times 10^{42}$</td>
<td>$1.4 \times 10^{42}$</td>
<td>$9.9 \times 10^{46}$</td>
<td>$1.4 \times 10^{-5}$</td>
<td>$37$</td>
<td>$3.8 \times 10^{-7}$</td>
<td>$2.1 \times 10^{-8}$</td>
<td>$2.0$</td>
</tr>
</tbody>
</table>


APPENDIX A: PROBABILITY DENSITY FUNCTIONS

We include figures showing the MCMC estimate of the marginalized probability distribution functions for each parameter in the SED modelling. These indicate whether a given parameter is well constrained or if it has only upper or lower limits. For example, the stellar population (SP) age for NGC 4594 is a lower limit but it is essentially unconstrained for NGC 4235, from inspection of plots A2 and A4. The summary statistics for the same are included in Tables 4 and 5.

This paper has been typeset from a \LaTeX/\BibTeX file prepared by the author.
Table 4. Best-fit Parameters from SED Modelling for NGC 4235

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<tr>
<th>Parameters</th>
<th>bestfit</th>
<th>median (−σ, +σ)</th>
<th>(−2σ, +2σ)</th>
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<td>Torus (Nenkova et al. 2008)</td>
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</tr>
<tr>
<td>σ</td>
<td>16.0</td>
<td>20.0 (−3.7, 8.4)</td>
<td>(−4.8, 22.5) deg</td>
</tr>
<tr>
<td>Y</td>
<td>20.3</td>
<td>19.2 (−5.4, 16.6)</td>
<td>(−9.1, 64.5)</td>
</tr>
<tr>
<td>( N_0 )</td>
<td>14.7</td>
<td>11.8 (−3.2, 2.2)</td>
<td>(−6.5, 3.0)</td>
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<tr>
<td>( q )</td>
<td>1.21</td>
<td>1.21 (−0.72, 0.72)</td>
<td>(−1.12, 1.18)</td>
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<tr>
<td>( \tau_c )</td>
<td>( 138.4 )</td>
<td>205.5 (−81.5, 65.8)</td>
<td>(−139.0, 90.0)</td>
</tr>
<tr>
<td>i</td>
<td>3.7</td>
<td>18.3 (−13.0, 18.5)</td>
<td>(−17.5, 36.4) deg</td>
</tr>
<tr>
<td>( C_{\text{trac}} )</td>
<td>47.2</td>
<td>55.2 (−8.9, 17.1)</td>
<td>(−12.1, 33.0) percent</td>
</tr>
<tr>
<td>hoa</td>
<td>62.1</td>
<td>56.6 (−13.4, 6.0)</td>
<td>(−30.8, 8.0) deg</td>
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<tr>
<td>esc Prob</td>
<td>100.00</td>
<td>99.98 (−9.20, 0.02)</td>
<td>(−67.11, 0.02) percent</td>
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<td>log Mass</td>
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<td>−0.49 (−0.24, 0.23)</td>
<td>(−0.50, 0.44) ( M_\odot )</td>
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<td>log Size</td>
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<td>−0.36 (−0.17, 0.26)</td>
<td>(−0.35, 0.61) pc</td>
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<td>ISM (Draine &amp; Li 2007)</td>
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<td>( u_{\text{min}} )</td>
<td>5.9</td>
<td>5.7 (−0.3, 0.3)</td>
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<tr>
<td>( q_{\text{PAH}} )</td>
<td>1.71</td>
<td>1.51 (−0.51, 0.73)</td>
<td>(−0.90, 1.90)</td>
</tr>
<tr>
<td>( f_{\text{PDR}} )</td>
<td>57.7</td>
<td>62.0 (−3.6, 3.5)</td>
<td>(−6.2, 8.6) percent</td>
</tr>
<tr>
<td>Stars (GRASIL; Silva et al. 1998)</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Age</td>
<td>12.9</td>
<td>8.4 (−5.0, 4.3)</td>
<td>(−7.8, 4.6) Gyr</td>
</tr>
<tr>
<td>Hot dust</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>( C_f )</td>
<td>33.2</td>
<td>54.1 (−29.3, 33.0)</td>
<td>(−46.1, 47.3) percent</td>
</tr>
<tr>
<td>Temp</td>
<td>1488.1</td>
<td>1398.5 (−211.5, 142.8)</td>
<td>(−435.6, 192.6) K</td>
</tr>
<tr>
<td>Amp</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>log AGNLum</td>
<td>42.541</td>
<td>42.500 (−0.145, 0.102)</td>
<td>(−0.264, 0.220) ergs/s</td>
</tr>
<tr>
<td>log ISMLum</td>
<td>42.714</td>
<td>42.714 (−0.007, 0.009)</td>
<td>(−0.012, 0.018) ergs/s</td>
</tr>
<tr>
<td>log SFR</td>
<td>−0.633</td>
<td>−0.633 (−0.007, 0.009)</td>
<td>(−0.012, 0.018) ( M_\odot/\text{yr} ) (Kennicutt 1998)</td>
</tr>
<tr>
<td>SL apcorr</td>
<td>0.233</td>
<td>0.225 (−0.064, 0.076)</td>
<td>(−0.132, 0.163)</td>
</tr>
</tbody>
</table>

**Torus:** The first six parameters determine the shape of the model torus spectrum. \( \sigma \) = angular scale height of the torus in degrees, \( Y \) = ratio of the outer radius to the inner radius of the torus set by sublimation, \( N_0 \) = average number of clouds along an equatorial sight-line, \( q \) = the power law index of the cloud distribution in radius - larger \( q \) values imply that clouds are more centrally concentrated, \( \tau_c \) = the optical depth at V-band (550 nm) through one cloud in the torus, \( i \) = the inclination of the torus relative to the sight-line, with 0 being face-on. The next five parameters are calculated from the torus parameters after the fit is done. \( C_{\text{trac}} \) = the fraction of sky covered in obscuring clouds as viewed from the centre of the torus, hoa = the half-opening angle: this is the polar angle at which the escape probability for accretion disk photons drops by 1/e relative to the poles, twice this angle is the apparent width of an ionisation cone if the torus were viewed edge on, esc Prob = escape probability of photons from the accretion disk or BLR for the best-fit sight-lines - this should be higher for type 1 AGN compared to type 2, log Mass = the mass of atomic + molecular hydrogen of the torus, log Size = an estimate of the outer radius of the torus assuming that the inner radius is set by 1500 K sublimation temperature and the best-fit values for \( Y \), ISM: \( u_{\text{min}} \) = a lower limit to the scaling of the interstellar radiation field (ISRF) relative to the Milky Way value, \( q_{\text{PAH}} \) = mass fraction of PAHs relative to total dust mass, \( f_{\text{PDR}} \) = fraction of dust luminosity from photo-dissociation regions (i.e. star-forming regions). **Stars:** Age = age of the model stellar population in Gyr. **Hot dust:** \( C_f \) and Temp are the sky covering fraction and dust temperature of the hot dust component. **Amp:** Amplitude of the torus component, log AGNLum = logarithm of the AGN bolometric luminosity, based on the scaling of the clumpy torus model needed for a best fit, log ISMLum = logarithm of the ISM luminosity, log SFR = logarithm of the star formation rate. SL apcorr: a measure of how much the large scale (LL) model spectrum of the diffuse ISM and stars has to be scaled down to fit the SL data. See Section 2.3 for details.
Figure 1. 610 MHz (top) and 325 MHz (bottom) contour images of NGC 4235. The contour levels are in percentage of the peak intensity and increase by factors of two. The peak intensity and lowest contour levels are $2.5 \times 10^{-2}$ Jy beam$^{-1}$ and ±0.3% for the top and $3.4 \times 10^{-1}$ Jy beam$^{-1}$ and ±0.074% for the bottom panel, respectively. The beam-size is $9.6 \times 5.9$ arcsec$^2$ at a PA of 53° at 610 MHz, and $10.6 \times 8.3$ arcsec$^2$ at a PA of 65° at 325 MHz. The grey scale range is 0–1.0 mJy beam$^{-1}$.
Figure 2. 325 – 610 MHz spectral index images in colour overlaid with 610 MHz radio contours for NGC 4235. The average spectral index is $\sim -0.6$ in the lobes and $\leq -1.8$ in the “relic” lobe. The contour levels are in percentage of the peak intensity and increase by factors of two. The peak intensity and lowest contour levels are $2.3 \times 10^{-2}$ Jy beam$^{-1}$ and $\pm 0.3\%$, respectively. The contour image is convolved with a circular beam of size 8 arcsec.

Table 5. Best-fit Parameters from SED Modelling for NGC 4594

<table>
<thead>
<tr>
<th>Parameters</th>
<th>bestfit</th>
<th>median ($-\sigma$, $+\sigma$)</th>
<th>($-2\sigma$, $+2\sigma$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Torus (Nenkova et al. 2008)</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\sigma$</td>
<td>23.0</td>
<td>25.3 ($-8.0$, 20.2)</td>
<td>$(-10.1$, 40.1) deg</td>
</tr>
<tr>
<td>$Y$</td>
<td>45.4</td>
<td>44.6 ($-31.7$, 37.5)</td>
<td>$(-38.7$, 52.8)</td>
</tr>
<tr>
<td>$N_0$</td>
<td>1.6</td>
<td>2.0 ($-0.8$, 2.8)</td>
<td>$(-1.0$, 9.1)</td>
</tr>
<tr>
<td>$q$</td>
<td>1.65</td>
<td>2.27 ($-0.85$, 0.52)</td>
<td>$(-1.89$, 0.70)</td>
</tr>
<tr>
<td>$\tau_v$</td>
<td>36.5</td>
<td>65.9 ($-48.3$, 148.0)</td>
<td>$(-55.1$, 220.9)</td>
</tr>
<tr>
<td>$i$</td>
<td>17.4</td>
<td>39.2 ($-26.9$, 30.5)</td>
<td>$(-37.4$, 47.5) deg</td>
</tr>
<tr>
<td>$C_{\text{frc}}$</td>
<td>33.0</td>
<td>43.4 ($-14.5$, 19.5)</td>
<td>$(-22.7$, 42.3) percent</td>
</tr>
<tr>
<td>hoa</td>
<td>69.1</td>
<td>62.9 ($-15.8$, 9.4)</td>
<td>$(-35.8$, 14.9) deg</td>
</tr>
<tr>
<td>esc Prob</td>
<td>99.99</td>
<td>89.53 ($-57.91$, 10.47)</td>
<td>$(-85.09$, 10.47) percent</td>
</tr>
<tr>
<td>log Mass</td>
<td>$-2.09$</td>
<td>$-2.27$ ($-0.52$, 0.54)</td>
<td>$(-0.94$, 1.26) M$\odot$</td>
</tr>
<tr>
<td>log Size</td>
<td>$-0.17$</td>
<td>$-0.22$ ($-0.54$, 0.29)</td>
<td>$(-0.93$, 0.48) pc</td>
</tr>
<tr>
<td>ISM (Draine &amp; Li 2007)</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$u_{\text{min}}$</td>
<td>3.9</td>
<td>3.6 ($-1.8$, 1.2)</td>
<td>$(-3.2$, 2.3)</td>
</tr>
<tr>
<td>$q_{\text{PAH}}$</td>
<td>0.51</td>
<td>0.63 ($-0.12$, 0.21)</td>
<td>$(-0.16$, 0.51)</td>
</tr>
<tr>
<td>$f_{\text{PDR}}$</td>
<td>17.5</td>
<td>14.8 ($-1.9$, 2.8)</td>
<td>$(-3.7$, 15.6) percent</td>
</tr>
<tr>
<td>Stars (GRASIL; Silva et al. 1998)</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Age</td>
<td>13.0</td>
<td>12.9 ($-0.2$, 0.1)</td>
<td>$(-0.4$, 0.1) Gyr</td>
</tr>
<tr>
<td>Amp</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>log AGNLum</td>
<td>42.148</td>
<td>42.077 ($-0.263$, 0.256)</td>
<td>$(-0.489$, 0.480) ergs/s</td>
</tr>
<tr>
<td>log ISMLum</td>
<td>42.455</td>
<td>42.470 ($-0.035$, 0.041)</td>
<td>$(-0.175$, 0.082) ergs/s</td>
</tr>
<tr>
<td>log SFR</td>
<td>$-0.892$</td>
<td>$-0.877$ ($-0.035$, 0.041)</td>
<td>$(-0.175$, 0.082) M$\odot$/yr (Kennicutt 1998)</td>
</tr>
</tbody>
</table>

See caption under Table 4.
Figure 3. 610 MHz (top) and 325 MHz (bottom) contour images of NGC 4594. The contour levels are in percentage of the peak intensity and increase by factors of two. The peak intensity and lowest contour levels are $6.1 \times 10^{-2}$ Jy beam$^{-1}$ and $\pm 0.15\%$ for the top and $2.1 \times 10^{-1}$ Jy beam$^{-1}$ and $\pm 0.15\%$ for the bottom panel, respectively. The beam-size is $8.8 \times 6.0$ arcsec$^2$ at a PA of 50$^\circ$ at 610 MHz, and $14.6 \times 8.4$ arcsec$^2$ at a PA of 45$^\circ$ at 325 MHz. The grey scale range is 0–1.4 mJy beam$^{-1}$. 
Figure 4. 325 – 610 MHz spectral index images in colour overlaid with 610 MHz radio contours for NGC 4594. The average spectral index is $+0.3$ for the southern lobe of NGC 4594. This value is however, likely to contaminated by galactic emission. The galaxy disk spectral index is $\sim -1.2$. The contour levels are in percentage of the peak intensity and increase by factors of two. The peak intensity and lowest contour levels are $6.7 \times 10^{-2}$ Jy beam$^{-1}$ and $\pm 0.15\%$, respectively. The contour image is convolved with a circular beams of size 10 arcsec.

Figure 5. The 610 MHz contour image of NGC 4594, showing the “spur” in the north-east direction. The contour levels are in percentage of the peak intensity and increase by factors of two. The peak intensity and lowest contour levels are $6.1 \times 10^{-2}$ Jy beam$^{-1}$ and $\pm 0.85\%$, respectively. The beam-size is $8.8 \times 6.0$ arcsec$^2$ at a PA of $50^\circ$. 

Figure 6. SED model-fitting of NGC 4235. The bestfit model is with a (i) clumpy torus, (ii) diffuse ISM, (iii) stellar population and (iv) a hot dust component. Blue (line, band) = (best-fit, acceptable-fit) Diffuse ISM component (Draine & Li models); Green = torus component (Nenkova et al. models); Orange = simple stellar population (large stellar population with short birth-time scale, allowed to age passively) (GRASIL models); Magenta = hot dust component; Red = total; Solid black line = IRS spectra; Black dots with triangles = 2MASS + IRAC measurements (as available); Open triangles = upper (pointed down) or lower (pointed up) limits, depending on aperture size relative to IRS. The residuals plot shows the range of acceptable models after the best-fit model has been subtracted. The dotted lines = ±10% residuals, roughly the expected systematic uncertainty based on the systematic uncertainty of the models and calibration.
Figure 7. SED model-fitting of NGC 4594. The bestfit model is with a (i) clumpy torus, (ii) diffuse ISM, and (iii) stellar population. For component details, see caption under Figure 6.
Figure A1. Marginalized posterior probability densities estimated from the MCMC samples for the torus parameters derived from the clumpyDREAM SED fits to NGC 4235.
Figure A2. Same as Figure A1, for the parameters of the stellar and ISM components.
Figure A3. Marginalized posterior probability densities estimated from the MCMC samples for the torus parameters derived from the clumpyDREAM SED fits to NGC 4594.
Figure A4. Same as Figure A3, for the parameters of the stellar and ISM components.